

Energy Dependence of the Analyzing Power for Proton-Carbon Scattering

A procedure based on Regge pole phenomenology is developed to determine the energy dependence of the analyzing power in pC elastic scattering in the CNI region. The model contains the Pomeron and two lower lying, $I=0$, Regge poles corresponding to the f_0 and the ω . It is shown that the measurement of the polarization P at one energy plus the measurement of the shape of the CNI peak at two energies is sufficient to predict the analyzing power at all energies. It is further shown that if the spin-flip factors for the f_0 and the ω are equal, the energy dependence can be predicted based on measurements of P and the shape at only one energy. Such information is available from E950. Using the Spin 2000 data from E950 we determine the real and imaginary parts of the spin-flip factor for $pL=21.7 \text{ GeV}/c$. From those and making the assumption just mentioned, we calculate the spin-flip factor for the Pomeron and for the lower lying Regge poles. With these the analyzing power at $pL=100$ and $250 \text{ GeV}/c$ is predicted. This simple assumption can be tested using the data recently obtained at $pL=100 \text{ GeV}/c$; if it fails that data can be used, without knowledge of P at $pL=100 \text{ GeV}/c$, to determine the three Regge spin-flip factors and thereby obtain the prediction of the model for the energy dependence, free from this assumption. The robustness of the assumption is tested by using two different Regge fits to the non-flip data. Comparison with the pp data from FermiLab E704 shows that the $I=1$ spin-flip factors must be much larger than the $I=0$ spin-flip factors. Considerable attention is given to the correlated errors at each stage.

Larry Trueman

Energy Dependence of CNI Analyzing Power
for proton-carbon scattering

$$P A_N(t) = \frac{N_p(t) - N_\downarrow(t)}{N_p(t) + N_\downarrow(t)}$$

small $|t| \leq 0.05 \text{ GeV}^2$

non-flip $f_0(s,t) = g_0(s,t) + g_0^{\text{em}}(s,t)$
 "nuclear" $\sim \frac{\alpha}{t}$

helicity flip $\cancel{f}_3(s,t) = \cancel{g}_3(s,t) + g_3^{\text{em}}(s,t)$
 $\sim \frac{\alpha k}{2m\sqrt{-t}}$

$$k = 1.79$$

$$A_{N(t)} \propto \text{Im}(f_0(s,t) f_3^*(s,t))$$

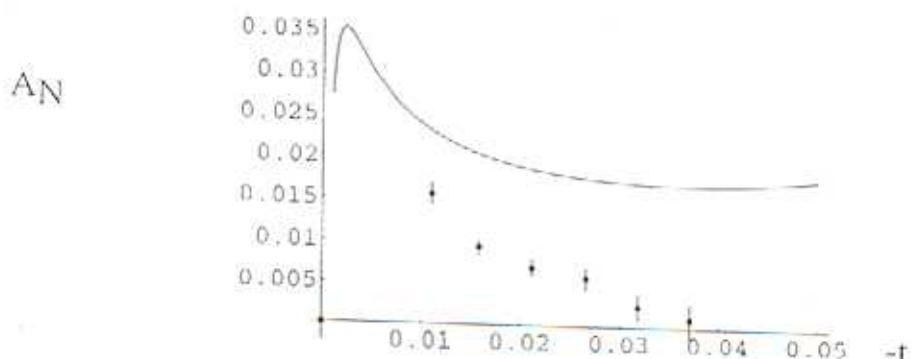
"pure" CNI $\text{Im}(g_0(s,t) g_3^{\text{em}}(s,t)) \propto \frac{\sigma_{tot}(s) e^{-kt}}{2m\sqrt{-t}}$

$\cancel{f}_3(s,t)$ but
 $\text{Im}(g_3(s,t) g_0^{\text{em}}(s,t))$ has same shape

parametrize $g_3(s,t) = \sqrt{-t} \frac{m}{M} g_0(s,t)$
 $= r_3 \frac{\sqrt{-t}}{M} \text{Im} g_0(s,t)$

$$\text{Im} r_3 = R_E t + \rho \text{Im} z, \quad R_E r_3 = -\text{Im} z + \rho R_E z$$

Pure CNI (no hadronic spin-flip, $\tau = 0$) compared with
Spin 2000 data.



To determine τ , fit ratio of measured asymmetry to pure CNI to

$$R(t) = \frac{A_N(t, \tau)}{A_N(t, 0)} = \left(1 - \frac{2}{\kappa} \text{Re}(\bar{\tau})\right) + \frac{2}{\kappa} \text{Im}(\bar{\tau}) \left(\frac{t}{t_c}\right) \frac{F^{(4)}}{F^{\text{empty}}} + \text{small}$$

$$t_c = -\frac{2\kappa \delta\sigma}{\sigma_{\text{tot}}} = -0.00127 \text{ GeV}^2$$

for $\pi^+ \pi^-$

If P is not known, can still determine shape parameter

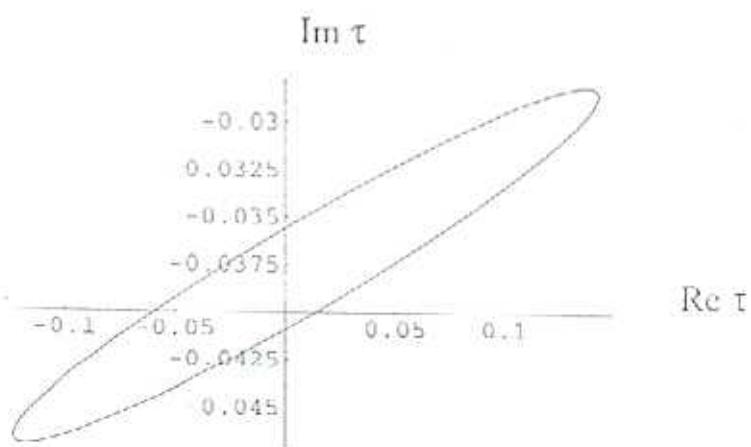
$$S(s) = \frac{\text{Im}T(s)}{1 - \frac{2}{\kappa} \text{Re}(s)}$$

Kopeliovich & Trueman:

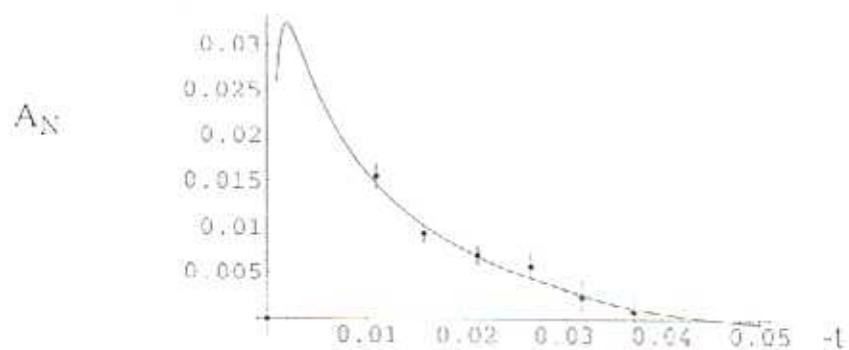
$$\tau_{pc} = \frac{1}{2} (\tau_{pp} + \tau_{pn})$$

(relationship nearly gone $T=0$)

68.3% confidence level ellipse for fit to Spin 2000 data



Best fit to Spin 2000 data for A_N
 $\tau = 0.010 - 0.038 i$



First try: Berger, Irving + Sorenson
 Regge pole expectations (1978)

7 Regge poles: Pomeron + 6 exchange-dep.
 Regge poles: 4 at $\alpha = \frac{1}{2}$
 2 at $\alpha = -\frac{1}{2}$
 for A_N

spin-dependence determined by "residues"

$$M_{\lambda_1 \lambda_2; \lambda_3}^i \propto S^{\alpha^{(4)}} \beta_{\lambda_1 \lambda_2}^{(4)} \beta_{\lambda_3}^{(4)}$$

Compares to Argonne data at $p_T = 6, 12 \text{ GeV}/c$

(Pomeron is assumed to be pole at $\alpha = 0$ so σ_{tot} is flat - not bad through pC fixed target range, but fails above)

(Comparisons are made for mainly above CERN region)

Use energy dependence of model to get

$$\bar{T}_{pp}(21.7) = 0.045 + 0.061 i$$

$$\bar{T}_{pC}(21.7) = 0.037 + 0.034 i \quad \text{n.b. p+inag part}$$

far from measured value $0.01 - 0.038 i$ at $21.7 \text{ GeV}/c$

(Consistent with large errors with $E704$
 $.03 \pm .02 i$)

Second try -

Plan -

Model : $g_0(s,t) = g_p(s) + g_f(s) + g_w(s)$

pp

$$g_p(s) = \chi \frac{(1+e^{-\pi s \alpha_p})}{s \sin \pi \alpha_p} s^{\alpha_p - 1}$$

$$g_f(s) = \gamma \frac{(1+e^{-\pi s \alpha_f})}{s \sin \pi \alpha_f} s^{\alpha_f - 1}$$

$$g_w(s) = \gamma' \frac{(1-e^{-\pi s \alpha_w})}{s \sin \pi \alpha_w} s^{\alpha_w - 1}$$

fit by Cudell et al

$$\alpha_p = 1.0933$$

$$\alpha_f = 0.64$$

$$\alpha_w = 0.44 \quad \text{etc.}$$

Assume

$$g_t(s,t) = \tau(s) \sqrt{\frac{-t}{m}} g_0(s,t)$$

$$= \sqrt{\frac{-t}{m}} \left\{ \tau_p g_p(s) + \tau_f g_f(s) + \tau_w g_w(s) \right\}$$

τ_p, τ_f, τ_w real constants

$$\tau_p = \frac{\beta_{+-}^p}{\beta_{++}^p} \quad \text{etc.}$$

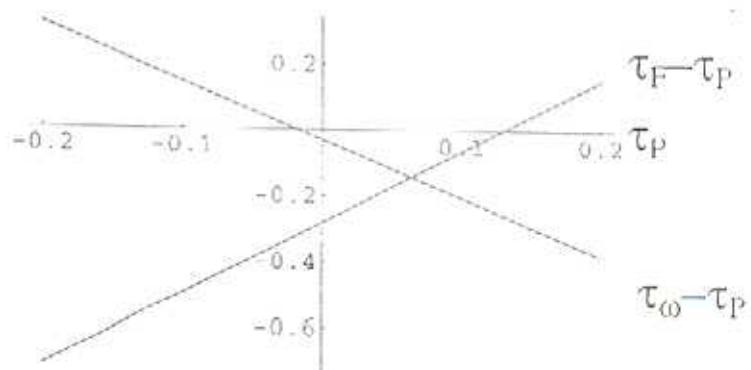
$$T(s) = \tau_p + (\tau_f - \tau_p) \frac{S_f(s)}{S_0(s)} + (\tau_w - \tau_p) \frac{g\omega(s)}{S_0(s)}$$

Knowing P at s_0 (E950) determine τ_e and τ_r of $T(s)$ since $(\tau_f - \tau_p)$ and $(\tau_w - \tau_p)$ are linear function of τ_p .

Then measure $S(s_1)$ without knowing P at s_1 and solve for τ_p .

Then A_N at all energies is known.

Solutions to equations generated by
measurement of $\tau(s)$ at $s=42$ with
known beam polarization



Speculative but reasonable assumption:

$$\bar{\tau}_w = \bar{\tau}_f = \bar{\tau}_R$$

allows determination of $\bar{\tau}_p$ and $\bar{\tau}_R$ from existing E950 data, and so ΔN at any energy can be calculated.

Get

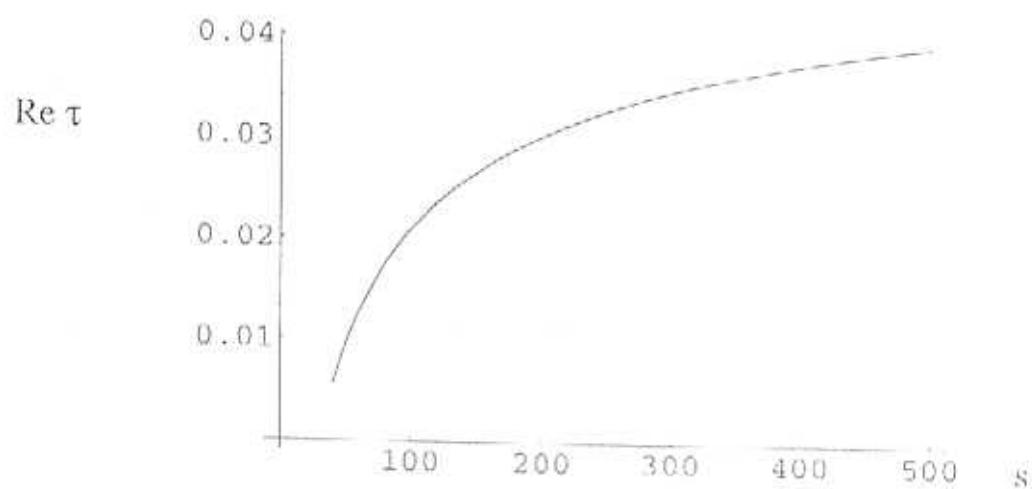
$$\bar{\tau}_p = 0.064$$

$$\bar{\tau}_R = -0.078$$

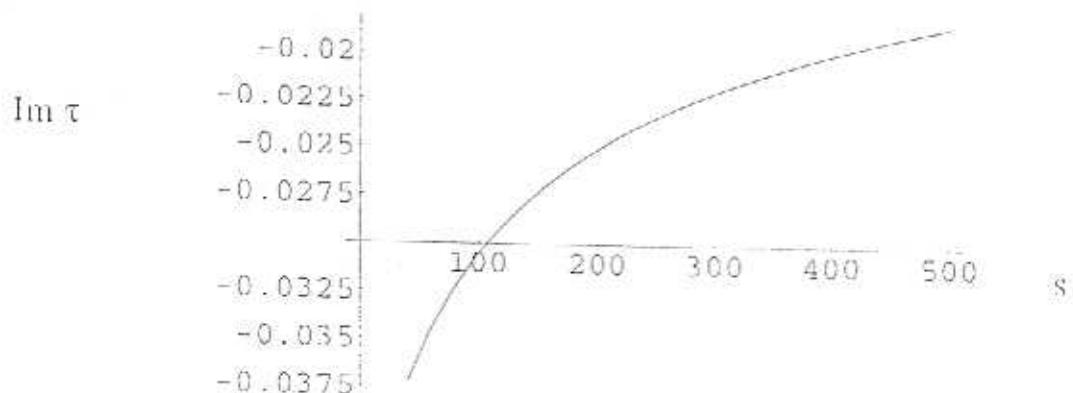
so $\text{Re } \bar{\tau}(s) \rightarrow 0.064$ as $s \rightarrow \infty$

$\text{Im } \bar{\tau}(s) \rightarrow 0$ from below

Energy dependence of τ

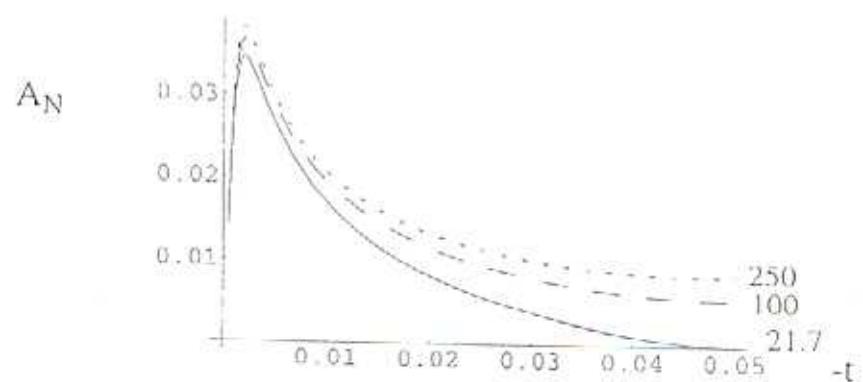


a.



b.

Analyzing power for $pL=100, 250$ GeV/c based on Regge model and fit at 21.7 GeV/c/



Stability & uniqueness?

Block et al (1998) use different form for Pomeron.

$$f_0^{PB} = i \left(A + B \left[\log \frac{S}{S_0} - i \frac{\pi}{2} \right]^2 \right)$$
$$+ i C_+^{\alpha-1} e^{i\pi(1-\alpha)/2} \quad (f, A_2)$$
$$+ \cancel{C_-}^{\alpha-1} s^{\alpha-1} e^{i\pi(1-\alpha)/2} \quad (\rho, \omega)$$

numerically Pomeron much more important relative to Regge than in Cadell et al
Both fits (σ_{tot} , ρ) equally good over our energy range.

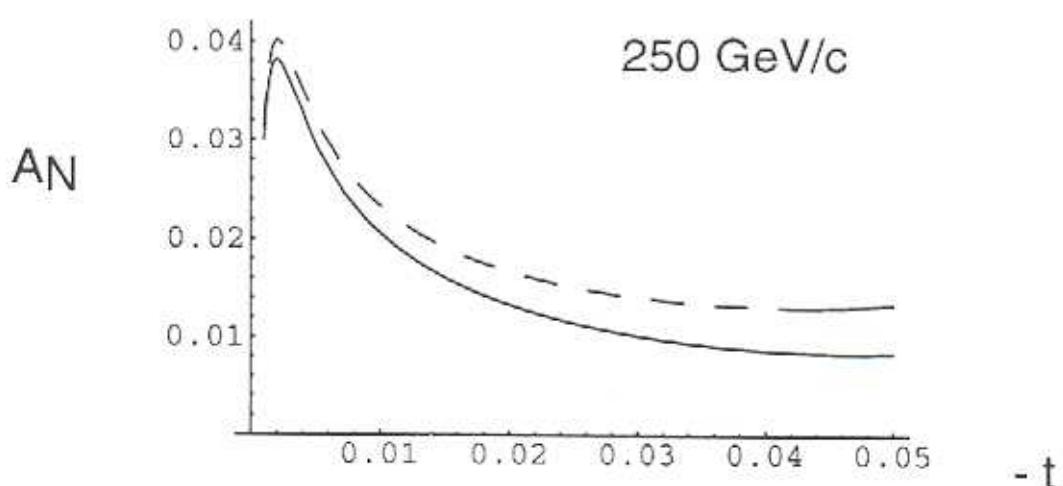
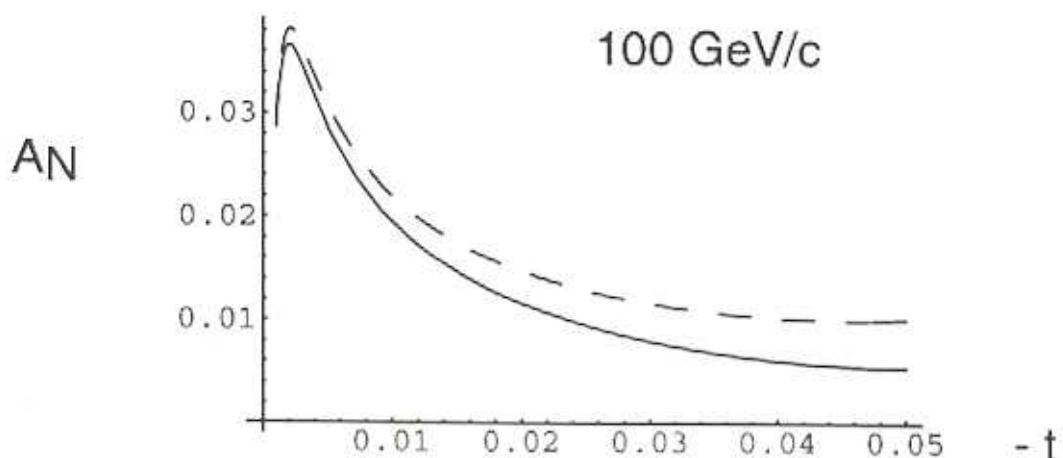
Results of same 2 pole game are somewhat different

P_t	Block et al	Cadell et al
100	$0.028 - 0.07i$	$0.033 - 0.026i$
250	$0.027 - 0.010i$	$0.043 - 0.019i$

1.5% confidence level ellipses don't overlap

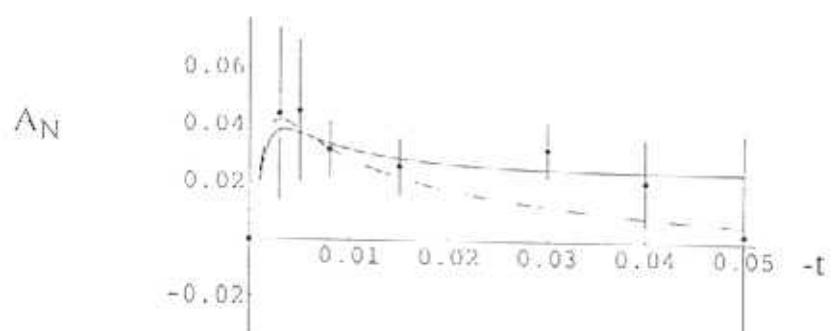
Difference about 5% near peak,
nearly a factor of 2. for $t = -0.05$

Comparison of projections for A_N to higher energy
using non-flip inputs of Cudell et al and of Block et al



solid line uses Cudell et al as input
dashed line uses Block et al as input

Λ_N for pp at $p_L = 200$ GeV/c with E704 low-t data. Solid curve is best fit with $\tau = 0.188 + .024 i$; dashed curves uses the values obtained from fit to Spin 2000 pC data extrapolated to $s=400$, $\tau = 0.041 - 0.02 i$.



Comparison of the error ellipse (68.3% confidence level) obtained from fit to E704 data below $|t|=0.05$ with that obtained by using Regge model applied to Spin 2000 data. Strong indication of large $I=1$ contribution to proton-proton scattering.

